A note on steady vortex flows in two dimensions

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Abstract

In this note, we give a general criterion for steady vortex flows in a planar bounded domain. More specifically, we show that if the stream function satisfies "locally" a semilinear elliptic equation with monotone or Lipschitz nonlinearity, then the flow must be steady.

Keywords: Euler equations, steady vortex flow, semilinear elliptic equation, stream function

1. Introduction and Main Result

Let $D \subset \mathbb{R}^2$ be a simply connected and bounded domain with a smooth boundary ∂D . The motion of an incompressible nonviscous fluid of unit density in D is governed by the following Euler dynamical equations:

$$\begin{cases} \partial_t \mathbf{v} + (\mathbf{v} \cdot \nabla) \mathbf{v} = -\nabla P, & (x, t) \in D \times \mathbb{R}^+, \\ \nabla \cdot \mathbf{v} = 0, & \\ \mathbf{v}(x, 0) = \mathbf{v}_0(x), & x \in D, \\ \mathbf{v} \cdot \mathbf{n} = 0, & (x, t) \in \partial D \times \mathbb{R}^+, \end{cases}$$
(1.1)

where $\mathbf{v} = (v_1, v_2)$ is the velocity field, P is the scalar pressure and **n** denotes the outward unit normal of ∂D . Here we impose the impermeability condition on the boundary.

Define the vorticity of the fluid as $\omega := \partial_1 v_2 - \partial_2 v_1$. Since **v** is divergence-free, there is a function ψ , called the stream function, such that $\mathbf{v} = J\nabla\psi := (\partial_2\psi, -\partial_1\psi)$, where J(a, b) = (b, -a) denotes clockwise rotation through $\frac{\pi}{2}$ for any planar vector (a, b). It is easy to see that ω and ψ satisfy the following Poisson's equation:

$$-\Delta \psi = \omega.$$

By the impermeability boundary condition, we deduce that ψ is a constant on ∂D . Without loss of generality, by adding a suitable constant we assume that ψ vanishes on the boundary.

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Thus ψ can be determined by ω in the following way

$$\psi(x,t) = G\omega(x,t) := \int_D G(x,y)\omega(y,t)dy, \ x \in D,$$

where G is the Green function for $-\Delta$ in D with zero Dirichlet boundary condition.

Taking the curl on both sides of the first equation in (1.1), we obtain

$$\partial_t \omega + J \nabla G \omega \cdot \nabla \omega = 0, \tag{1.2}$$

which is a nonlinear transport equation for ω and is usually called the vorticity equation.

In this paper, we are concerned with the steady vorticity equation, that is,

$$J\nabla G\omega \cdot \nabla \omega = 0. \tag{1.3}$$

To motivate the definition of weak solutions to the steady vorticity equation, we multiply any $\phi \in C_c^{\infty}(D)$ on both sides of (1.3) and integrate by parts formally to obtain

$$\int_{D} \omega J \nabla G \omega \cdot \nabla \phi dx = 0.$$
(1.4)

It is not difficult to check that the integral in (1.4) makes sense if $\omega \in L^{4/3}(D)$. In fact, for $\omega \in L^{\frac{4}{3}}(D)$, by L^p estimate we have $G\omega \in W^{2,\frac{4}{3}}(D)$, thus $G\omega \in W^{1,4}(D)$ by Sobolev embedding, therefore the integral in (1.4) makes sense by Hölder's inequality.

Definition 1.1. We call $\omega \in L^{4/3}(D)$ a weak solution to the steady vorticity equation (1.3) if it satisfies

$$\int_{D} \omega J \nabla G \omega \cdot \nabla \phi dx = 0, \ \forall \phi \in C_{c}^{\infty}(D).$$
(1.5)

In the past several decades, various methods have been proposed to construct steady vortex flows. The most commonly used method is to investigate the following semilinear elliptic problem

$$\begin{cases} -\Delta \psi = f(\psi), & x \in D, \\ \psi = 0, & x \in \partial D, \end{cases}$$
(1.6)

where $f : \mathbb{R} \to \mathbb{R}$ is a Lipschitz function. It is obvious that

$$J\nabla\psi\cdot\nabla f(\psi) = J\nabla\psi\cdot(f'(\psi)\nabla\psi) = 0$$
 a.e. in D ,

which means that any solution to (1.6) corresponds to a steady vortex flow in classical sense with ψ as the stream function. See [6][8][15][14] and the references listed therein. Another efficient way to construct steady vortex flows is called the vorticity method, which was first established by Arnold [1][2]. See also [3][4][9][10][11][12][16]. Roughly speaking, the vorticity method states that any steady vortex flow is equivalent to a critical point of the kinetic energy subject to

some appropriate constraints for the vorticity. For example, Turkington in [16] considered the maximization of the kinetic energy

$$E(\omega) := \frac{1}{2} \int_D \int_D G(x, y) \omega(x) \omega(x) dx dy$$

in the admissible class

$$K_{\lambda}(D) := \{ \omega \in L^{\infty}(D) \mid 0 \le \omega \le \lambda \text{ a.e. in } D, \ \int_{D} \omega(x) dx = 1 \}.$$
(1.7)

Turkington proved the existence of a maximizer and showed that any maximizer ω^{λ} must be of the form

$$\omega^{\lambda} = \lambda I_{A^{\lambda}}, \ A^{\lambda} = \{ x \in D \mid G\omega^{\lambda}(x) > \mu^{\lambda} \},$$
(1.8)

where $I_{A^{\lambda}}$ denotes the characteristic function of A^{λ} , and μ^{λ} is a Lagrange multiplier depending on λ . To show that ω^{λ} is a steady solution to the vorticity equation (1.3), one can not use the form (1.8) anymore, since the nonlinearity here is a Heaviside function with discontinuity at μ^{λ} and the regularity of ∂A^{λ} is unknown(still an open question). To show that ω^{λ} satisfies (1.5), Turkington used the fact that ω^{λ} is an energy maximizer in $K_{\lambda}(D)$. See [16] for the detailed proof.

In [5], Burton proved that if ω belongs to $L^{4/3}(D)$ and satisfies $\omega = f(G\omega)$ a.e. in D, where $f : \mathbb{R} \to \mathbb{R} \cup \{\pm \infty\}$ is a monotone function, then ω is a weak solution to (1.3). By Burton's result, in order to obtain a steady vortex flow from (1.8), we do not need additional information about the energy of ω^{λ} .

Another example of steady vortex flows with vorticity concentrated in multiple separated regions is given in [7]. Therein the authors studied the following elliptic problem

$$\begin{cases} -\Delta \psi = \lambda \sum_{i=1}^{k} I_{A_i}, \ x \in D, \\ A_i = B_{\delta}(x_{0,i}) \cap \{ x \in D \mid \psi(x) > \kappa_i \}, \end{cases}$$
(1.9)

where I_{A_i} denotes the characteristic function of A_i , λ is a given positive number, κ_i is a real number depending on λ , $x_{0,i}$ is a given point in D, δ is a very small positive number such that $B_{\delta}(x_{0,i}) \subset D$ and $\overline{B_{\delta}(x_{0,i})} \cap \overline{B_{\delta}(x_{0,j})} = \emptyset$ if $i \neq j$. The authors constructed a solution to (1.9) for sufficiently large λ , such that each A_i is a simply connected domain bounded by a C^1 closed curve and is strictly contained in $B_{\delta}(x_{0,i})$ (or equivalently, $dist(A_i, \partial B_{\delta}(x_{0,i})) > 0$). To show that ψ satisfies (1.5), one can integrate by parts directly since each ∂A_i is C^1 and ψ is continuous across ∂A_i .

Notice that in (1.9) the vorticity $\omega = -\Delta \psi$ is no longer a function of the stream function ψ , since the k Lagrange multipliers $\kappa_i, \dots, \kappa_k$ may be different numbers. However, the vorticity is a function of the stream function "locally".

Our aim in this note is to give a general criterion for solutions of the steady vorticity equation (1.3), that is, if the stream function satisfies "locally" a semilinear elliptic equation with monotone or Lipschitz nonlinearity, then the corresponding flow must be steady.

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Before stating the theorem, we give some notations for clarity. We will use supp(f) to denote the support of some function f, and the distance between two planar sets A and B is defined by

$$dist(A,B) := \inf_{x \in A, y \in B} |x - y|.$$

Let $\delta > 0$ be a positive number, the notation A_{δ} denotes the δ -neighbourhood in D of some planar set A, or equivalently,

$$A_{\delta} := \{ x \in D \mid dist(x, A) < \delta \}.$$

Theorem 1.2. Let k be a positive integer. Suppose that $\omega \in L^{4/3}(D)$ satisfies

$$\omega = \sum_{i=1}^{k} \omega_i, \ \min_{1 \le i < j \le k} \{ dist(supp(\omega_i), supp(\omega_j)) \} > 0, \ \omega_i = f^i(G\omega), a.e. \ in \ supp(\omega_i)_{\delta}$$
(1.10)

for some $\delta > 0$, where each f^i is either monotone from \mathbb{R} to $\mathbb{R} \cup \{\pm \infty\}$ or Lipschitz continuous from \mathbb{R} to \mathbb{R} , then ω is a weak solution to the steady vorticity equation (1.3).

Remark 1.3. In the above theorem, if f_i is Lipschitz continuous from \mathbb{R} to \mathbb{R} , then ω_i must be bounded since $G\omega \in L^{\infty}(D)$ by L^p estimate and Sobolev embedding.

Remark 1.4. Examples of steady vortex flows satisfying (1.10) can also be found in [12], where each $f^i : \mathbb{R} \to \mathbb{R} \cup \{\pm \infty\}$ is an unknown nondecreasing function.

2. Proof of the Main Result

In this section, we give the proof of Theorem 1.2. The basic idea is to approximate each f^i by a sequence of bounded Lipschitz functions.

For a function $f : \mathbb{R} \to \mathbb{R} \cup \{\pm \infty\}$, we shall use the following notation in the rest of this paper for convenience:

 $C_f := \{ s \in \mathbb{R} \mid f \text{ is continuous at } s \},$ $D_f := \{ s \in \mathbb{R} \mid f \text{ is not continuous at } s \}.$

Lemma 2.1. Suppose that $f : \mathbb{R} \to \mathbb{R} \cup \{\pm \infty\}$ is a monotone function, then there exists a sequence of bounded and smooth functions $\{f_n\}$ such that

$$|f_n(s)| \le |f(s)|, \ \forall s \in \mathbb{R},$$
$$\lim_{n \to +\infty} f_n(s) = f(s), \ \forall s \in C_f.$$

Proof. Without loss of generality we assume that f is nondecreasing and bounded (we can use truncation if f is unbounded).

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First we consider the case f is nonnegative. Let ρ be the standard mollifier of one dimension, that is,

$$\rho(s) = \begin{cases} c_0 e^{-\frac{1}{1-s^2}}, & |s| < 1, \\ 0, & |s| \ge 1, \end{cases}$$

where c_0 is a positive number such that $\int_{\mathbb{R}} \rho(s) ds = 1$. Note that $\rho \in C_c^{\infty}(\mathbb{R})$. Define

$$f^{\varepsilon}(s) = \int_{-\infty}^{+\infty} \rho_{\varepsilon}(s-\varepsilon-r)f(r)dr = \int_{s-2\varepsilon}^{s} \rho_{\varepsilon}(s-\varepsilon-r)f(r)dr,$$

where $\rho_{\varepsilon}(s) := \varepsilon^{-1} \rho(s \varepsilon^{-1})$ with $\varepsilon > 0$ as a parameter. It is easy to check that $f^{\varepsilon} \in C^{\infty}(\mathbb{R})$. Since f is nonnegative and nondecreasing, we have

$$|f^{\varepsilon}(s)| = \left| \int_{s-2\varepsilon}^{s} \rho_{\varepsilon}(s-\varepsilon-r)f(r)dr \right| \le \left| \int_{s-2\varepsilon}^{s} \rho_{\varepsilon}(s-\varepsilon-r)f(s)dr \right| = |f(s)|, \ \forall s \in \mathbb{R}.$$

Moreover, for any $s \in C_f$,

$$|f^{\varepsilon}(s) - f(s)| = \left| \int_{s-2\varepsilon}^{s} \rho_{\varepsilon}(s - \varepsilon - r)f(r)dr - f(s) \right|$$
$$= \left| \int_{s-2\varepsilon}^{s} \rho_{\varepsilon}(s - \varepsilon - r)(f(r) - f(s))dr \right|$$
$$\leq \sup_{r \in [s-2\varepsilon,s]} |f(r) - f(s)|$$

which goes to 0 as $\varepsilon \to 0^+$. Thus we have proved the lemma for nonnegative f.

For the case f is non-positive, we can define

$$f^{\varepsilon}(s) = \int_{-\infty}^{+\infty} \rho_{\varepsilon}(s+\varepsilon-r)f(r)dr = \int_{s}^{s+2\varepsilon} \rho_{\varepsilon}(s+\varepsilon-r)f(r)dr,$$

then by repeating the above argument we can prove the lemma for non-positive f.

When f is a general nondecreasing function, we write $f = f^+ - f^-$, where $f^+(s) := \max\{f(s), 0\}$ and $f^-(s) := -\min\{f(s), 0\}$. According to the above discussion, we can choose two sequences of smooth functions $\{f_n^+\}$ and $\{f_n^-\}$ such that

$$|f_n^+(s)| \le |f^+(s)|, \ |f_n^-(s)| \le |f^-(s)|, \ \forall s \in \mathbb{R},$$
$$\lim_{n \to +\infty} f_n^+(s) = f^+(s), \ \lim_{n \to +\infty} f_n^-(s) = f^-(s), \ \forall s \in C_f.$$

Here we used the fact that f^+ and f^- are both continuous on C_f . The lemma is proved by choosing $f_n = f_n^+ - f_n^-$.

Proof of Theorem 1.2: For $i = 1, \dots, k$, if f^i is a monotone function, by Lemma 2.1 we can choose a sequence of bounded and smooth functions $\{f_n^i\}$ such that

$$|f_n^i(s)| \le |f^i(s)|, \ \forall s \in \mathbb{R},$$

$$(2.1)$$

$$\lim_{n \to +\infty} f_n^i(s) = f^i(s), \ \forall s \in C_f.$$
(2.2)

If f^i is Lipschitz continuous, we can also choose a sequence of bounded Lipschitz functions $\{f_n^i\}$ satisfying (2.1)(2.2) by using truncation.

Since $\omega \in L^{4/3}(D)$, by L^p estimate we have $G\omega \in W^{2,4/3}(D)$, then by Sobolev embedding we obtain $G\omega \in W^{1,4}(D)$. By the chain rule for Sobolev functions(see [13], 4.22), it is easy to verify that

$$J\nabla G\omega \cdot \nabla (f_n^i(G\omega)) = (f_n^i)'(G\omega)J\nabla G\omega \cdot \nabla G\omega = 0 \text{ for a.e. } x \in supp(\omega_i)_{\delta},$$
(2.3)

where we used the fact $J\nabla G\omega \cdot \nabla G\omega \equiv 0$. Since $|f_n^i(s)| \leq |f^i(s)|$ for each $s \in \mathbb{R}$ and $n = 1, 2, \cdots$, we deduce that

$$supp(f_n^i(G\omega)) \cap supp(\omega_i)_{\delta} \subset supp(\omega_i).$$

Define $\omega_n = \sum_{i=1}^k f_n^i(G\omega) I_{supp(\omega_i)_{\delta}}$. Taking into account (1.10) and (2.3), we can easily check that ω_n belongs to $W^{1,4}(D)$ and satisfies

$$|\omega_n| \le \sum_{i=1}^k |f_n^i(G\omega)I_{supp(\omega_i)_{\delta}}| \le \sum_{i=1}^k |f^i(G\omega)I_{supp(\omega_i)_{\delta}}| = |\omega| \quad \text{a.e. } x \in D,$$
(2.4)

$$J\nabla G\omega \cdot \nabla \omega_n = 0 \text{ a.e. in } D.$$
(2.5)

Therefore we obtain

$$\int_{D} \omega_n J \nabla G \omega \cdot \nabla \phi dx = 0, \ \forall \phi \in C_c^{\infty}(D).$$
(2.6)

Now we claim that

$$\lim_{n \to +\infty} \omega_n = \omega \text{ a.e. in } D.$$
(2.7)

In fact, it suffices to show that for each i

$$\lim_{n \to +\infty} f_n^i(G\omega(x)) = f^i(G\omega(x)) \text{ for a.e. } x \in supp(\omega_i)_{\delta}.$$

For $x \in (G\omega)^{-1}(C_{f^i})$, by (2.2) we have $f_n^i(G\omega(x)) \to f^i(G\omega(x))$. So we need just consider the case $x \in (G\omega)^{-1}(D_{f^i})$. Since each f^i is either monotone or Lipschitz continuous, the set D_{f^i} is countable, thus it suffices to show that for each $s \in D_{f^i}$ there holds

$$\lim_{n \to +\infty} f_n^i(G\omega(x)) = f^i(G\omega(x)) \text{ for a.e. } x \in (G\omega)^{-1}(s)$$

To show this, first we use the property of Sobolev functions (see [13], 4.22) to obtain

$$\omega = -\Delta G \omega = 0$$
 a.e. on $(G \omega)^{-1}(s)$,

then by (2.4) we have

$$\omega_n = 0$$
 a.e. on $(G\omega)^{-1}(s)$,

therefore

$$\lim_{n \to +\infty} \omega_n = \omega \text{ a.e. on } (G\omega)^{-1}(s).$$

Combining (2.4),(2.6) and (2.7), we are able to apply the dominated convergence theorem to obtain

$$\int_D \omega J \nabla G \omega \cdot \nabla \phi dx = 0,$$

which is the desired result.

Remark 2.2. According to the proof of Theorem 1.2, we need only impose the following two abstract conditions on f^i :

- (1) D_{f^i} is a countable set;
- (2) there exist a sequence of bounded Lipschitz functions $\{f_n^i\}$ and a constant C > 0 such that

$$|f_n^i(s)| \le C|f^i(s)|, \ \forall s \in \mathbb{R},$$
$$\lim_{n \to +\infty} f_n^i(s) = f^i(s), \ \forall s \in C_f.$$

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